

$\forall u_0 \in H, \exists! u(t, u_0)$ soln in $([0, T], H), \forall T$

$$P_t \phi(u_0) = \mathbb{E}(\phi(u(t, u_0)))$$

$(\nabla P_t) = P_t^* v = v$ \exists follows from tightness of $\frac{1}{t} \int_0^t u_s ds$

Uniqueness (same case)

$du = (Au + u^3) dt + \phi dw$, Dirichlet bdy
 ϕ s.t. $\exists!$ global soln ($\phi = I$)

2 invariant measures ν_1, ν_2

u_0^1, u_0^2 for ν_1, ν_2

$$\rightarrow u^1, u^2$$

$$r = u^1 - u^2$$

$$\frac{dr}{dt} = AR - (u^1)^3 - (u^2)^3$$

$$\frac{1}{2} \frac{d|r|^2}{dt} + |r_x|^2 \leq 0$$

$$\frac{1}{2} \frac{d|r|^2}{dt} + \lambda |r|^2 \leq 0$$

$$\underline{\mathbb{E}|r|^2} \leq e^{-2\lambda t} \underline{\mathbb{E}|u_0^1 - u_0^2|^2}$$

$$\mathbb{E}(\phi(u^1(t))) - \mathbb{E}(\phi(u^2(t))) \rightarrow 0$$

$$\downarrow \int \phi d\nu_1 \quad \quad \quad \int \phi d\nu_2$$

if ϕ Lipschitz

$$\rightarrow \nu_1 = \nu_2$$

Kolmogorov eqn: $du = (Au + f(u))dt + \phi dw$

f : Burgers, R.P or N.S, or

$$Y(t, u_0) = \mathbb{E}(\phi(u(t, u_0))) = P_t \phi(u_0)$$

Ito Formula: $d\phi(u(t, u_0)) = (Au + f(u), D\phi(u(t, u_0))) + (D\mu\phi(u(t, u_0)), \phi dW) + \frac{1}{2} \text{Tr} \phi \phi^* D^2 \phi(u(t, u_0))$

Take $E \rightarrow t=0$ and then t_0

Kol. eqn $\rightarrow \begin{cases} \frac{d\psi}{dt} = \frac{1}{2} \text{Tr} \phi \phi^* D^2 \psi + (Au + f(u), D\phi) = L\psi \\ \psi(0) = \phi \end{cases}$

Much more than for $\phi=0$ because of $\text{Tr} \phi \phi^* D^2 \psi$

• $\phi = I: \text{Tr} D^2 \psi = \Delta \psi$

\rightarrow uniqueness is easier.

• ϕ degenerates $\rightarrow \text{rank } \phi < \infty$

\rightarrow hypoelliptic case

In finite dimension, the adjoint of the K eq.

\rightarrow Fokker-Planck eq.

In infinite dimension:

\rightarrow no Lebesgue measure

gradient system:

\int_{recall}

ex: reaction diffusion: $\rightarrow v(u) = \frac{1}{4} \int |u|^4 d\xi, Dv(u) = u^3$

$du = (Au + Dv(u)) dt + dW$

$(A = \Delta + \text{Dir bdy cond})$

$L\psi = \frac{1}{2} \text{Tr} D^2 \psi + (Au + Dv(u), D\psi)$

$\frac{1}{2} \int L\psi e^{-v(u)} dv(u)$

$v = N(0, \frac{1}{2} (-A)^{-1})$

inv measure for $v=0$

$\frac{1}{2} \int \text{Tr} D^2 \psi e^{-v(u)} e^{-\frac{1}{2} (Au)^2} du = - (D\psi(u), v'(u)) - (D\psi(u), Au)$

$$\int L \phi e^{-V(\omega)} d\nu = 0$$

$$\frac{d}{dt} \int P_t \phi(\omega) e^{-V(\omega)} d\nu = \int \rho d\mu \quad \mu = \frac{1}{Z} e^{-V}$$

$$\left. \begin{aligned} &= \int L P_t \phi(\omega) e^{-V(\omega)} d\nu \\ &= 0 \end{aligned} \right\} \Rightarrow \mu = \frac{1}{Z} e^{-V} \text{ invariant measure}$$

6) Rough, but not too rough, noise

Reaction diffusion with space-time white noise

$$du = (\Delta u + u^3 - u) dt + dW$$

Periodic bdy condition on $[0, 1]^d$

$d=1$ easy

$d=2$ critical $dz = Az + dW \quad z \in C([0, T], H^s) \quad \forall s > 0$

$d=3$: rough \rightarrow MH

In fact $z \in C([0, T], B_{p,r}^{s'})$

$$s < 0, p, r \geq 1$$

$$u = \gamma + z \quad A = \Delta - I \quad dz = Az + dW$$

$$d\gamma/dt = A\gamma + (\gamma + z)^3 = A\gamma + \gamma^3 + 3\gamma^2 z + 3\gamma z^2 + z^3$$

z is gaussian: $z \sim N(0, \frac{1}{2}(-A)^{-1})$

$$z \notin A, \text{Tr}(-A)^{-1} = \infty$$

Up to renormalization

z^p can be defined

- where this renormalized z^p lives?

- is γ sufficiently smooth to define γz^2 and γz^3

Gaussian calculus, Wick products

Facts: $\int_{H^{-1}} (x, (-A)^{1/2} y)^2 d\mu(x) = |y|_H^2 = \mathbb{E} \left(z, (-A)^{1/2} y \right)^2$

$$x \rightsquigarrow (x, (-A)^{1/2} y)$$

exists in $L^2(H^{-1}, \mu)$ for $y \in H$

$$\int e^{(x, (-A)^{1/2} y)} d\mu(x) = e^{1/2 |y|_H^2}$$

Hermite polynomials.

$$e^{-t^2/2 + t\xi} = \sum_0^{\infty} \frac{t^m}{m!} H_m(\xi)$$

$$\int H_m((x, (-A)^{1/2} y)) H_n((x, (-A)^{1/2} \bar{y})) d\mu = \delta_{n,m} (y, \bar{y})^m$$

$$V_n = \text{Sp} \left\{ x \rightsquigarrow H_n(x, (-A)^{1/2} y), |y|=1 \right\}$$

$$L^2(H^{-1}, \mu) = \bigoplus_{n \in \mathbb{N}} V_n$$

nth Wiener class

$$z^3(\xi) \quad \xi \in (0,1)$$

$$z^3(\xi) = \lim_{\varepsilon \rightarrow 0} (z, \rho_\varepsilon(\cdot - \xi))^3$$

$$\rightarrow x \rightsquigarrow (x, (-A)^{1/2} y)^3$$

with $y = (-A)^{-1/2} \rho_\varepsilon(\cdot - \xi)$

$$\Pi_3 \left(x \rightsquigarrow (x, (-A)^{1/2} y)^3 \right) = \sqrt{3} H_3 \left(x \rightsquigarrow (x, (-A)^{1/2} y)^3 \right)$$

↑
proj onto V_3 for $|y|=1$

$$z_\varepsilon^3(\xi) = (z, \rho_\varepsilon(\cdot - \xi)) = K_\varepsilon^3(z, (-A)^{1/2} \tilde{\rho}_{\varepsilon, \xi})$$

$$\tilde{\rho}_{\varepsilon, \xi} = \frac{1}{K_\varepsilon} (-A)^{-1/2} \rho_\varepsilon(\cdot - \xi)$$

$$K_\varepsilon = \|(-A)^{-1/2} \rho_\varepsilon(\cdot - \xi)\|_H \xrightarrow{\varepsilon \rightarrow 0} \infty$$

$$\pi_3(z_\varepsilon(\xi)^3) = K_\varepsilon^3 \sqrt{\Gamma} \cdot H_3\left(\frac{z_\varepsilon(\xi)}{K_\varepsilon}\right) = z_\varepsilon(\xi)^3 \underbrace{K_\varepsilon^{-3}}_{\text{renormalization}}$$

If z has low μ then
 $\pi_3(z_\varepsilon(\xi))$ converges

$\pi_3(z_\varepsilon^3)$ converges in $L^2(\Omega; H^s)$ $\forall s < 0$
 and in $L^k(\Omega, B_{p,r}^s)$ $\forall s < 0$
 ↖ Besov space

$$\rightarrow :z^3 := \lim_{\varepsilon \rightarrow 0} \pi_3(z_\varepsilon^3) \in L^k(\Omega, B_{p,r}^s) \quad \forall s < 0, k, p, r \geq 1$$

Something for $:z^2$:

$$dy/dt = Ay + y^3 + 3y^2z + 3yz^2 + z^3$$

$$\rightarrow dy/dt = Ay + y^3 + 3y^2z + 3yz^2 + :z^3:$$

Look for $y \in C([0, T], B_{p,r}^s) \cap L^k(0, T, B_{p,r}^s)$

for well chosen k, p , and p large

$$\bullet (v^1, v^2) \rightsquigarrow v^1 v^2$$

$$B_{p,r}^a \times B_{p,r}^b \rightarrow B_{p,r}^{a+b-2/r}$$

$$a+b > 0 \text{ if } a, b < 2/p$$

$$\bullet t \rightsquigarrow :z(t)^3: \in C([0, T], B_{p,r}^s) \cap L^k(0, T, B_{p,r}^s) = C_{p,r,k,s} T$$

$$:z^3: \in L^k(0, T, B_{p,r}^s)$$

$$\rightarrow \forall u_0 \in B_{p,r}^s, \exists! u \in C([0, T], B_{p,r}^s) \\ T = T^*(u)$$

$$\nu(dx) := \frac{1}{2} e^{-\int :x^4: d\bar{x}} \mu(dx) \text{ invariant measure}$$

→ global existence for ν a.e. u_0

Mourrant-Weber - proved global existence for every u_0
(not just with inv measure)

• First argument of step 0 of MH.